

# From an initial data to a global solution of the nonlinear Schrödinger equation: a building process

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## Abstract

The purpose of this work is to construct a continuous map from the homogeneous Besov space  $\dot{B}_{2,4}^0(\mathbb{R}^2)$  in the set  $\mathcal{G}$  of initial data in  $\dot{B}_{2,4}^0(\mathbb{R}^2)$  which gives birth to global solution of the mass critical non linear Schrödinger equation in the space  $L^4(\mathbb{R}^{1+2})$ . We use the fact that solutions of scale which are different enough almost do not interact; the main point is that we determine a condition about the size of the scale which depends continuously on the data.

## 1 Introduction

We consider in the space  $\mathbb{R}^2$  the non linear Schrödinger equation of the type

$$(GNLS_m) \quad \begin{cases} i\partial_t u + \Delta u &= P_3(\bar{u}, u) \\ u|_{t=0} &= u_0(x) \end{cases} \quad (1)$$

where  $P_3$  is a homogeneous polynomial of order 3 in  $z$  and  $\bar{z}$ .

Let us observe that if we define  $u_\lambda(t, x) = \lambda u(\lambda^2 t, \lambda x)$  for a solution of  $(GNLS_m)$  then  $u_\lambda$  is also a solution. As we work in  $\mathbb{R}^2$ , the  $L^2$  norm, which represents the total mass, is invariant under this scaling transformation. This family of equations is thus called "mass critical".

As done for instance in [4], Strichartz estimate for the Schrödinger equation claims that for any couples  $(p_j, q_j)$  such that

$$\frac{2}{p_j} + \frac{2}{q_j} = 1 \quad \text{and} \quad (p_j, q_j) \neq (2, \infty),$$

there exists a constant  $C$  such that for any interval  $I = [I^-, I^+]$  of  $\mathbb{R}$ , one has

$$\|u\|_{L^{p_1}(I; L^{q_1}(\mathbb{R}^2))} \lesssim \|u(I^-)\|_{L^2(\mathbb{R}^2)} + \|i\partial_t + \Delta u\|_{L^{p_2}'(I; L^{q_2}'(\mathbb{R}^2))}. \quad (2)$$

It is now classical to prove, with a fixed point argument, that Equation  $(GNLS_m)$  is locally well posed for initial data in  $L^2(\mathbb{R}^2)$ , and globally well posed for small initial data in  $L^2(\mathbb{R}^2)$ . More precisely, we have the following theorem (see for instance [4]).

**Theorem 1.1.** *Let  $u_0$  be an initial data in  $L^2(\mathbb{R}^2)$ . Then a positive time  $T$  exists such that there is a unique solution to  $(GNLS_m)$  in  $L^4([0, T] \times \mathbb{R}^2)$ . Moreover, a positive constant  $c_0$  exists such that, if  $\|u_0\|_{\dot{B}_{2,4}^0(\mathbb{R}^2)}$  is less than or equal to  $c_0$ , the solution is global and belongs to  $L^4(\mathbb{R}^{1+2})$  and*

$$\int_{\mathbb{R}^{1+2}} |u(t, x)|^4 dx dt \lesssim \|u_0\|_{L^2(\mathbb{R}^2)}^4. \quad (3)$$

Let us mention that some important particular cases of  $(GNLS_m)$  have been intensively studied: the case when  $P_3(u, \bar{u}) = -|u|^2 u$ , called defocusing case and the case when  $P_3(u, \bar{u}) = |u|^2 u$  called the focusing case. In the defocusing case, B. Dodson proves in [7] that for any initial data in  $L^2(\mathbb{R}^2)$ , a unique global solution exists in  $L^4(\mathbb{R}^{1+2})$ ; he proved in [8] that it was also the case in the defocusing case when the mass of the initial data was less than the mass of the ground state. In the works [11] and [10], refined blow phenomena are described for initial data closed to the ground state.

In the present work, we proved general results far away from using the refined structure of the focusing or defocusing case. As observed by P. Gérard in an published work (see the text [13] by F. Planchon and references therein), it is possible to improve a little this result using Besov spaces. For technical reasons, we prefer to use a continuous version of the Littlewood-Paley theory. Let us consider a radial function  $\varphi$  in  $\mathcal{D}(\mathbb{R}^d \setminus \{0\})$  such that

$$\varphi(\xi) = f(|\xi|) \quad \text{with} \quad \int_0^\infty f(r^{-1}) \frac{dr}{r} = 1.$$

The fact that the measure  $\mu^{-1} d\mu$  is invariant under dilatation gives, for any  $\xi$  in  $\mathbb{R}^d \setminus \{0\}$ ,

$$\int_0^\infty \varphi(\mu^{-1}\xi) \frac{d\mu}{\mu} = 1.$$

Let us define  $\Delta_\mu a \stackrel{\text{def}}{=} \mathcal{F}^{-1}(\varphi(\mu^{-1}\xi)\widehat{a}(\xi))$ . Now let us recall the definition of Besov spaces.

**Definition 1.1.** We define the Besov space as the space of distributions such that

$$\|u\|_{\dot{B}_{p,r}^s(\mathbb{R}^d)} \stackrel{\text{def}}{=} \|\mu^s \|\Delta_\mu u\|_{L^p}\|_{L^r(\mathbb{R}^+; \frac{d\mu}{\mu})}.$$

The theory is strictly analogous to the one of the discrete case. In this context, the norm  $\dot{B}_{2,4}^0(\mathbb{R}^2)$  is defined by

$$\|u\|_{\dot{B}_{2,4}^0(\mathbb{R}^2)}^4 \stackrel{\text{def}}{=} \int_0^\infty \|\Delta_\mu u\|_{L^2(\mathbb{R}^2)}^4 \frac{d\mu}{\mu}.$$

Let us observe that the space  $\dot{B}_{2,4}^0(\mathbb{R}^2)$  is slightly bigger than  $L^2(\mathbb{R}^2) = \dot{B}_{2,2}^0(\mathbb{R}^2)$ , but has the same scaling invariant. We have the following result (see [13]).

**Theorem 1.2.** Let  $u_0$  be an initial data in  $\dot{B}_{2,4}^0(\mathbb{R}^2)$ . Then a positive time  $T$  exists such that there is a unique solution to  $(GNLS_m)$  in  $L^4([0, T] \times \mathbb{R}^2)$ . Moreover, a positive constant  $c_0$  exists such that, if  $\|u_0\|_{\dot{B}_{2,4}^0(\mathbb{R}^2)}$  is less than or equal to  $c_0$ , the solution is global and belongs to  $L^4(\mathbb{R}^{1+2})$  and

$$\int_{\mathbb{R}^{1+2}} |u(t, x)|^4 dx dt \lesssim \|u_0\|_{\dot{B}_{2,4}^0}^4. \quad (4)$$

Let us give some sketchy indications about the proof of this theorem. Using an improvement of Strichartz inequality (2) proved by J. Bourgain in [2], we get

$$\|e^{it\Delta} u_0\|_{L^4(\mathbb{R}^{1+2})} \lesssim \|u_0\|_{\dot{B}_{2,4}^0(\mathbb{R}^2)}. \quad (5)$$

Then it is possible to make a fixed point theorem in the space  $L^4([0, T] \times \mathbb{R}^2)$ . Let us point out that if  $u$  is a solution to  $(GNLS_m)$  in the space  $L^4([0, T] \times \mathbb{R}^2)$  then  $w = u - e^{it\Delta} u_0$  satisfies

$$(i\partial_t w + \Delta w) = P_3(u, \bar{u}) \in L^{\frac{4}{3}}([0, T] \times \mathbb{R}^2) \quad \text{and} \quad w|_{t=0} = 0.$$

Strichartz estimate (2) with  $(p_2, q_2) = (4, 4)$  and  $(p_1, q_1) = (\infty, 2)$  implies that  $u - e^{it\Delta} u_0$  belongs to the space  $L^\infty([0, T]; L^2(\mathbb{R}^2))$ .

**Definition 1.2.** Let us denote by  $\mathcal{G}$  the set of all initial data in  $\dot{B}_{2,4}^0(\mathbb{R}^2)$  such that a global solution of  $(GNLS_m)$  exists in  $L^4(\mathbb{R}^{1+2})$ .

It is classical that the set  $\mathcal{G}$  is an open subset of  $\dot{B}_{2,4}^0(\mathbb{R}^2)$ . Indeed, if  $u_0$  belongs to  $\mathcal{G}$ , and  $v_0$  to  $\dot{B}_{2,4}^0(\mathbb{R}^2)$ , one can search the solution associated with  $u_0 + v_0$  under the form  $u + v$ , where  $u$  is the solution associated with  $u_0$ . Elementary computations lead to

$$i\partial_t v + \Delta v = E(u, v) \quad \text{with} \quad |E(u, v)(t, x)| \lesssim |v(t, x)|(|u(t, x)|^2 + |v(t, x)|^2).$$

Inequalities (5) and (2) allow to prove that if  $v_0$  is small enough in the norm  $\dot{B}_{2,4}^0(\mathbb{R}^2)$ , then fixed point argument works in the space  $L^4(\mathbb{R}^{1+2})$  for  $v$ .

The purpose of this work is to prove the following theorem.

**Theorem 1.3.** *A continuous function  $F$  from  $\dot{B}_{2,4}^0(\mathbb{R}^2)$  into  $\mathcal{G}$ , and a constant  $C_0$  exist, such that*

$$C_0^{-1} \|v\|_{\dot{B}_{2,4}^0(\mathbb{R}^2)} \leq \|F(v)\|_{\dot{B}_{2,4}^0(\mathbb{R}^2)} \leq C_0 \|v\|_{\dot{B}_{2,4}^0(\mathbb{R}^2)}.$$

As far as we know, there is no results of this type in the literature. The motivation of this result is to catch some informations about the set  $\mathcal{G}$  in the spirit of the works [5] by J.-Y. Chemin and I. Gallagher or [6] by J.-Y. Chemin, I. Gallagher and P. Zhang in the context of homogeneous incompressible Navier-Stokes equation.

## 2 Structure of the proof

The idea is to make the function  $v$  small by multiplication, and then to copy, as many times as required, a dilatation of the (now small) function. Let us introduce, for  $p$  in  $[1, \infty[$ , the following notation:

$$\forall v \in L^2(\mathbb{R}^d), \quad \pi_{\mathbb{S}}^{(p)} v \stackrel{\text{def}}{=} \frac{v}{\|v\|_{L^p}} \quad \text{and} \quad \pi_{\mathbb{S}}^B v = \frac{v}{\|v\|_{\dot{B}_{2,4}^0}}.$$

The first step consists in the definition of a function  $\mathcal{F}$  on  $\dot{B}_{2,4}^0(\mathbb{R}^2) \times \mathbb{R}_+^* \times \mathbb{R}_+$ , with values in  $\dot{B}_{2,4}^0(\mathbb{R}^2)$ , which make explicit the idea explained above.

**Definition 2.1.** *Let us consider a smooth function  $\chi$  defined on  $\mathbb{R}^+$  such that*

$$\chi|_{[0, \frac{1}{4}]} = 0 \quad \text{and} \quad \chi|_{[\frac{3}{4}, +\infty[} = 1.$$

Let us define  $\mathcal{F}$  from  $\dot{B}_{2,4}^0(\mathbb{R}^2) \times \mathbb{R}_+^* \times \mathbb{R}_+$  into  $\dot{B}_{2,4}^0(\mathbb{R}^2)$  as follows

$$\mathcal{F}(v, \Lambda, \lambda) = \sum_{j=0}^{[\lambda]} \Lambda^{-j} v(\Lambda^{-j} \cdot) + \chi(\lambda - [\lambda]) \Lambda^{-[\lambda]-1} v(\Lambda^{-[\lambda]-1} \cdot).$$

The main properties of  $\mathcal{F}$  are described hereafter.

**Proposition 2.1.** *The map  $\mathcal{F}$  is continuous from  $\dot{B}_{2,4}^0(\mathbb{R}^2) \times \mathbb{R}_+^* \times \mathbb{R}_+$  into  $\dot{B}_{2,4}^0(\mathbb{R}^2)$ . Moreover, there exists a positive constant  $C$  and a continuous function  $\Lambda_0$  from  $]1, +\infty[ \times \mathbb{S}(\dot{B}_{2,4}^0(\mathbb{R}^2))$  into  $\mathbb{R}_+$  such that, for any  $v$  in  $\dot{B}_{2,4}^0(\mathbb{R}^2)$  different from 0, we have*

$$\Lambda \geq \Lambda_0(\lambda + 1, \pi_{\mathbb{S}}^B v) \implies \frac{1}{C}(\lambda + 1) \|v\|_{\dot{B}_{2,4}^0(\mathbb{R}^2)}^4 \leq \|\mathcal{F}(v, \Lambda, \lambda)\|_{\dot{B}_{2,4}^0(\mathbb{R}^2)}^4 \leq C(\lambda + 1) \|v\|_{\dot{B}_{2,4}^0(\mathbb{R}^2)}^4. \quad (6)$$

This proposition means that functions, the scales of which are different enough, are, in some sense, "orthogonal" in the Besov space  $\dot{B}_{2,4}^0(\mathbb{R}^2)$ . The main point of the proposition is that the choice of the scales size can be made continuously with respect to the function  $v$ .

**Proposition 2.2.** *Let  $B_{c_0}$  be the open centered ball, the radius  $c_0$  of which is given by Theorem 1.2. There exists a continuous function  $\Lambda_1$  from  $]1, \infty[ \times B_{c_0}$  into  $]0, +\infty[$  such that for any  $v$  in  $B_{c_0}$ ,*

$$\Lambda \geq \Lambda_1(\lambda + 1, v) \implies \mathcal{F}(v, \Lambda, \lambda) \in \mathcal{G}.$$

The proof of this proposition relies on the idea that two solutions of an evolution equations with scales that are different enough almost do not interact. The idea which is now currently used goes back to the works of H. Bahouri and P. Gérard (see [1]) and F. Merle and L. Vega (see [12]). Again, it is related to the fact that the choice of the size can be made continuously with respect to the function  $v$ .

These two propositions imply Theorem 1.3. Indeed, if  $C_0$  is the constant given by Theorem 1.2, we define

$$F(v) \stackrel{\text{def}}{=} \left(1 - \chi(\rho^{-1}\|v\|_{\dot{B}_{2,4}^0(\mathbb{R}^2)})\right)v + \chi(\rho^{-1}\|v\|_{\dot{B}_{2,4}^0(\mathbb{R}^2)})\mathcal{F}(\rho\pi_{\mathbb{S}}^B v, \Lambda(v), \|v\|_{\dot{B}_{2,4}^0(\mathbb{R}^2)}) \quad \text{with}$$

$$\Lambda(v) \stackrel{\text{def}}{=} \Lambda_0(\|v\|_{\dot{B}_{2,4}^0(\mathbb{R}^2)}^4 + 1, \pi_{\mathbb{S}}^B v) + \Lambda_1(\|v\|_{\dot{B}_{2,4}^0(\mathbb{R}^2)}^4 + 1, \rho\pi_{\mathbb{S}}^B v).$$

and then apply the above two propositions.

The idea of the proof of Proposition 2.1 is elementary: we develop the quantity  $\|\mathcal{F}(v, \Lambda, \lambda)\|_{\dot{B}_{2,4}^0(\mathbb{R}^2)}^4$  and we prove that the cross terms are small enough thanks to the following lemma.

**Lemma 2.3.** *Let us denote by  $\mathbb{S}(L^p(\mathbb{R}^d))$  the unit sphere of the Lebesgue space  $L^p(\mathbb{R}^d)$ . For any  $p$  in  $]1, +\infty[$ , there exists a continuous function  $\Lambda_p$  from  $]0, 1[ \times \mathbb{S}(L^p(\mathbb{R}^d)) \times \mathbb{S}(L^{p'}(\mathbb{R}^d))$  such that for any couple  $(f, g)$  of  $(L^p(\mathbb{R}^d) \setminus \{0\}) \times (L^{p'}(\mathbb{R}^d) \setminus \{0\})$ , one has*

$$\Lambda \geq \Lambda_p(\varepsilon, \pi_{\mathbb{S}}^{(p)} f, \pi_{\mathbb{S}}^{(p')} g) \implies \left| \int_{\mathbb{R}^d} f(x) \Lambda^{-\frac{d}{p'}} g(\Lambda^{-1}x) dx \right| \leq \varepsilon \|f\|_{L^p(\mathbb{R}^d)} \|g\|_{L^{p'}(\mathbb{R}^d)}.$$

If one does not take continuity into account, the property is a very classical one: it is enough to approach, up to  $\varepsilon$ , the functions  $\pi_{\mathbb{S}}^{(p)} f$  and  $\pi_{\mathbb{S}}^{(p')} g$  respectively in  $L^p$  and  $L^{p'}$ , by ring-supported functions. Then, one chooses  $\Lambda$  large enough such that the integral of the truncated functions is 0. The question is: can the choice of the cut off be made continuously with respect to the functions? The answer is no, as shown by the following picture:

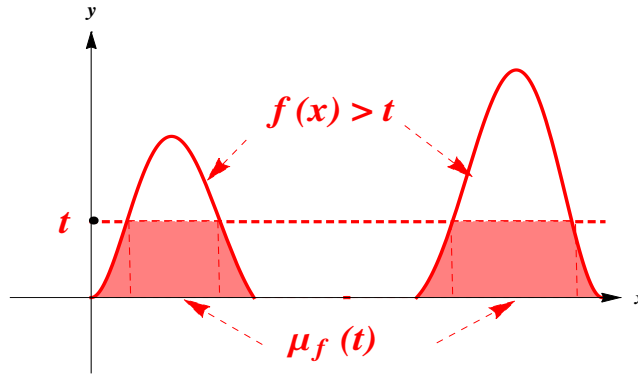


Figure 1: The graph of a non-negative function and measurable function  $f$

If the function is radial and non increasing, it becomes possible. The use of the non-increasing reordering enables one to prove Lemma 2.3.

Section 3 is, first, devoted to brief recalls on the definition and main properties of the symmetric non-increasing rearrangement. One proves, then, Lemma 2.3, which enables one to continuously truncate the rearrangements. Once it is done, one can prove Proposition 2.1.

Section 4 concentrates on the proof of Proposition 2.2, where Lemma 2.3 plays a major role regarding continuity.

Section 5 is devoted to the proof of basic results on the non-increasing reordering which are collected here for the reader's convenience.

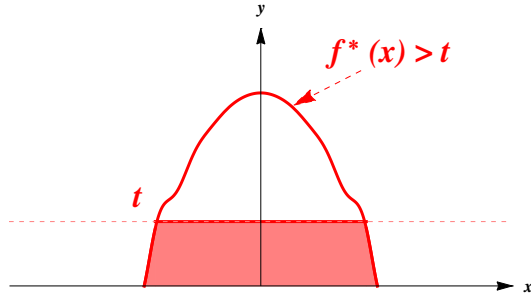


Figure 2: The graph of the rearrangement  $f^*$

### 3 Non-increasing reordering and almost-orthogonality

To begin with, let us work in the space  $\mathbb{R}^d$ . Let us, first, recall basic notations, that we will require later on. For any measurable function  $f$ , we set

$$(f > t) \stackrel{\text{def}}{=} \{x \in \mathbb{R}^d / f(x) > t\} \quad \text{and} \quad (f \leq t) \stackrel{\text{def}}{=} \{x \in \mathbb{R}^d / f(x) \leq t\}.$$

Let us consider a non-negative function and measurable function  $f: \mathbb{R}^d \rightarrow \mathbb{R}$ , which vanishes at infinity, in the sense that, for any strictly positive real number  $t$ , the set  $(|f| > t)$  is of finite measure.

**Definition 3.1.** Let  $A$  be a measurable set of  $\mathbb{R}^d$ , the Lebesgue measure of which is finite.  $A^*$  denotes the open centered ball whose measure is the same of the one of  $A$ . Let us consider a real-valued, measurable function  $f$ , defined on  $\mathbb{R}^d$ , which vanishes at infinity, in the sense that, for any strictly positive real number  $t$ , the set  $(|f| > t)$  is of finite measure. The symmetric non-increasing rearrangement of  $f$  is the measurable and non negative function  $f^*$ , defined on  $[0, +\infty[$  by

$$f^*(x) = \int_0^{+\infty} \mathbb{1}_{(|f|>t)^*(x)} dt.$$

More detailed information can be found, for instance, in [9] and [16]. One can already notice that  $f^* = |f|^*$ , and that the function  $f^*$  is a radially symmetric and non-increasing one, in the sense that  $f^*(x) = f^*(y)$  if  $|x| = |y|$ , and  $f^*(x) \geq f^*(y)$  if  $|x| \leq |y|$ .

We will require the non-increasing reordering properties described hereafter.

**Proposition 3.1.** Let us consider two real-valued, measurable functions  $f$  and  $g$ , defined on  $\mathbb{R}^d$ , which vanishes at infinity. For any strictly positive number  $s$ ,

$$(f(s \cdot))^* = f^*(s \cdot), \tag{7}$$

$$\int_{\mathbb{R}^d} f(x)g(x)dx \leq \int_{\mathbb{R}^d} f^*(x)g^*(x)dx \quad \text{and} \tag{8}$$

$$\int_{\mathbb{R}^d} f^*(x) \mathbb{1}_{(g^* \leq s)}(x) dx \leq \int_{\mathbb{R}^d} f(x) \mathbb{1}_{(g \leq s)}(x) dx. \tag{9}$$

If, moreover,  $f$  and  $g$  belong to the space  $L^p$ , it is also the case of  $f^*$  and  $g^*$ ; in addition,

$$\|f\|_{L^p} = \|f^*\|_{L^p} \quad \text{and} \quad \|f^* - g^*\|_{L^p} \leq \|f - g\|_{L^p}.$$

For the convenience of the reader, this proposition is proved in Section 5.

*Proof of Lemma 2.3* Thanks to inequalities (7) and (8), one is led to find an upper bound to

$$\left| \int_{\mathbb{R}^d} f(x) \Lambda^{-\frac{d}{p'}} g(\Lambda^{-1}x) dx \right| \leq I_\Lambda \stackrel{\text{def}}{=} \int_{\mathbb{R}^d} f^*(x) \Lambda^{-\frac{d}{p'}} g^*(\Lambda^{-1}x) dx.$$

As explained in the introduction, this means truncating the functions  $f^*$  and  $g^*$  in rings, which can be done using the following lemma.

**Lemma 3.2.** For any  $p$  in  $[1, \infty[$ , there exists a continuous function  $R_0$  from  $]0, 1[ \times \mathbb{S}(L^p)$  to  $\mathbb{R}^+$  such that, for any  $f$  in  $\mathbb{S}(L^p)$ ,

$$\|\phi_{R_0(\varepsilon, f)} f^*\|_{L^p} = \varepsilon \quad \text{with} \quad \phi_R(x) \stackrel{\text{def}}{=} 1 - \mathbb{1}_{\{R^{-1} \leq |x| \leq R\}}(x).$$

*Proof.* Let us first prove that, for any positive valued function  $h$ , radially symmetric, non-increasing, the  $L^p$  norm of which is equal to 1, and for any real number  $\varepsilon$  in  $]0, 1[$ , there exists a unique  $R(\varepsilon, h)$  such that

$$\int_{\mathbb{R}^d} \phi_{R(\varepsilon, h)}(x) h^p(x) dx = \varepsilon.$$

From the dominated convergence theorem, the function

$$\rho_h(R) \stackrel{\text{def}}{=} \int_{\mathbb{R}^d} \phi_R(x) h^p(x) dx$$

is continuous and non-increasing on  $[1, \infty[$ . Moreover :  $\rho_h(1) = \|h\|_{L^p}^p = 1$ , and  $\rho_h(+\infty) = 0$ . Let us denote by  $\underline{R}_h$  the maximum of numbers  $R$  satisfying  $\rho_h(R) = 1$ . The function  $\rho_h$  is decreasing on the interval  $[\underline{R}_h, \infty[$ . Thus, there exists a unique  $R_h$  such that  $\rho_h(R_h) = \varepsilon$ . Applying this property to the function  $h = f^*$  enables one to define the function  $R_0$ .

In order to prove the continuity of  $R_0$ , let us consider a sequence  $(\varepsilon_n, f_n)$  of  $]0, 1[ \times \mathbb{S}(L^p)$ , which converges towards  $(\varepsilon, f)$  in  $]0, 1[ \times \mathbb{S}(L^p)$ . It is enough to prove that any convergent subsequence converges towards  $R_0(\varepsilon, f)$ . Let us denote by  $R_\infty$  the limit of a convergent subsequence of  $R_0(\varepsilon_n, f_n)$ . For sake of simplicity, extraction will not be distinguished. The triangle inequality leads to

$$\|\phi_{R_0(\varepsilon_n, f_n)} f_n^*\|_{L^p} - \|f^* - f_n^*\|_{L^p} \leq \|\phi_{R_0(\varepsilon_n, f_n)} f^*\|_{L^p} \leq \|\phi_{R_0(\varepsilon_n, f_n)} f_n^*\|_{L^p} + \|f^* - f_n^*\|_{L^p}.$$

By definition of  $R_0$ ,

$$\varepsilon_n - \|f^* - f_n^*\|_{L^p} \leq \|\phi_{R_0(\varepsilon_n, f_n)} f^*\|_{L^p} \leq \varepsilon_n + \|f^* - f_n^*\|_{L^p}.$$

Since the symmetric non-increasing rearrangement is 1-lipschitzian on the Lebesgue spaces  $L^p$ , one has

$$\varepsilon_n - \|f - f_n\|_{L^p} \leq \|\phi_{R_0(\varepsilon_n, f_n)} f^*\|_{L^p} \leq \varepsilon_n + \|f - f_n\|_{L^p}. \quad (10)$$

The Lebesgue theorem implies then that:

$$\lim_{n \rightarrow +\infty} \|\phi_{R_0(\varepsilon_n, f_n)} f^*\|_{L^p} = \|\phi_{R_\infty} f^*\|_{L^p}.$$

By passing through the limit in (10), we obtain  $\|\phi_{R_\infty} f^*\|_{L^p(\mathbb{R}^d)} = \varepsilon$ , which, by definition of  $R_0$ , leads to  $R_\infty = R_0(\varepsilon, f)$ . Then, the function  $R_0$  is a continuous one. The lemma is proved.  $\square$

*Conclusion of the proof of Lemma 2.3* One can estimate  $I_\Lambda$ . Thanks to the identities

$$f^* = \phi_{R_0(\varepsilon, f)} f^* + \mathbb{1}_{\frac{1}{R_0(\varepsilon, f)} \leq |x| \leq R_0(\varepsilon, f)} f^* \quad \text{and} \quad g^* = \phi_{R_0(\varepsilon, g)} g^* + \mathbb{1}_{\frac{1}{R_0(\varepsilon, g)} \leq |x| \leq R_0(\varepsilon, g)} g^*,$$

we have

$$\begin{aligned} I_\Lambda &\leq \int_{\mathbb{R}^d} (\phi_{R_0(\varepsilon, f)} f^*)(x) \Lambda^{-\frac{d}{p'}} g^*\left(\frac{x}{\Lambda}\right) dx \\ &\quad + \int_{\mathbb{R}^d} (\mathbb{1}_{\frac{1}{R_0(\varepsilon, f)} \leq |x| \leq R_0(\varepsilon, f)} f^*)(x) \Lambda^{-\frac{d}{p'}} (\phi_{R_0(\varepsilon, g)} g^*)\left(\frac{x}{\Lambda}\right) dx \\ &\quad + \int_{\mathbb{R}^d} (\mathbb{1}_{\frac{1}{R_0(\varepsilon, f)} \leq |x| \leq R_0(\varepsilon, f)} f^*)(x) \Lambda^{-\frac{d}{p'}} (\mathbb{1}_{\frac{1}{R_0(\varepsilon, g)} \leq |x| \leq R_0(\varepsilon, g)} g^*)\left(\frac{x}{\Lambda}\right) dx \\ &\leq \|\phi_{R_0(\varepsilon, f)} f^*\|_{L^p} \|g\|_{L^{p'}} + \|f^*\|_{L^p} \|\phi_{R_0(\varepsilon, g)} g^*\|_{L^{p'}} \\ &\quad + \int_{\mathbb{R}^d} (\mathbb{1}_{\frac{1}{R_0(\varepsilon, f)} \leq |x| \leq R_0(\varepsilon, f)} f^*)(x) \Lambda^{-\frac{d}{p'}} (\mathbb{1}_{\frac{1}{R_0(\varepsilon, g)} \leq |x| \leq R_0(\varepsilon, g)} g^*)\left(\frac{x}{\Lambda}\right) dx. \end{aligned}$$

As  $\|f\|_{L^p} = \|f^*\|_{L^p} = \|g\|_{L^{p'}} = \|g^*\|_{L^{p'}} = 1$ , this leads to

$$\begin{aligned} I_\Lambda &\leq 2\varepsilon + \int_{\mathbb{R}^d} \left( \mathbb{1}_{\frac{1}{R_0}(\varepsilon, f) \leq |x| \leq R_0(\varepsilon, f)} f^*(x) \Lambda^{-\frac{d}{p'}} \left( \mathbb{1}_{\frac{1}{R_0(\varepsilon, g)} \leq |x| \leq R_0(\varepsilon, g)} g^*(x) \left( \frac{x}{\Lambda} \right) \right) dx \right. \\ &\leq 2\varepsilon + \int_{\mathbb{R}^d} \mathbb{1}_{\frac{1}{R_0(\varepsilon, f)} \leq |x| \leq R_0(\varepsilon, f)}(x) \mathbb{1}_{\frac{\Lambda}{R_0(\varepsilon, g)} \leq |x| \leq \Lambda R_0(\varepsilon, g)}(x) f^*(x) \Lambda^{-\frac{d}{p'}} g^*\left(\frac{x}{\Lambda}\right) dx. \end{aligned}$$

Let us notice that if

$$\Lambda \geq 2R_0(\varepsilon, f) R_0(\varepsilon, g) \quad (11)$$

then

$$\left\{ x \in \mathbb{R}^d, \frac{1}{R_0(\varepsilon, f)} \leq |x| \leq R_0(\varepsilon, f) \right\} \cap \left\{ x \in \mathbb{R}^d, \frac{\Lambda}{R_0(\varepsilon, g)} \leq |x| \leq \Lambda R_0(\varepsilon, g) \right\} = \emptyset.$$

Thus, provided the condition (11) is satisfied, one has  $I_\Lambda(f, g) \leq 2\varepsilon$ . For  $(f, g)$  in  $\mathbb{S}(L^p) \times \mathbb{S}(L^{p'})$ , let us define

$$\Lambda_p(\varepsilon, f, g) \stackrel{\text{def}}{=} 2R_0\left(\frac{\varepsilon}{2}, f\right) R_0\left(\frac{\varepsilon}{2}, g\right).$$

This function is suitable and Lemma 2.3 is thus proved.  $\square$

*Proof of Proposition 2.1* A basic expansion and the use of the scaling invariance lead to

$$\begin{aligned} &\left| \|\mathcal{F}(v, \Lambda, \lambda)\|_{\dot{B}_{2,4}^0}^4 - (([\lambda] + 1) + \chi(\lambda - [\lambda])) \|v\|_{\dot{B}_{2,4}^0(\mathbb{R}^2)}^4 \right| \lesssim \mathcal{E}_{v, \Lambda, \lambda} \quad \text{where} \\ \mathcal{E}_{v, \Lambda, \lambda} &\stackrel{\text{def}}{=} \int_0^\infty \sum_{(j_1, j_2, j_3, j_4) \in \tilde{\mathcal{J}}} \|\Delta_\mu(\Lambda^{-j_1} v(\Lambda^{-j_1} \cdot))\|_{L^2} \|\Delta_\mu(\Lambda^{-j_2} v(\Lambda^{-j_2} \cdot))\|_{L^2} \\ &\quad \times \|\Delta_\mu(\Lambda^{-j_3} v(\Lambda^{-j_3} \cdot))\|_{L^2} \|\Delta_\mu(\Lambda^{-j_4} v(\Lambda^{-j_4} \cdot))\|_{L^2} \frac{d\mu}{\mu} \quad \text{with} \\ \tilde{\mathcal{J}} &\stackrel{\text{def}}{=} \{0, \dots, [\lambda] + 1\}^4 \setminus \{(j, j, j, j), j \in \{0, \dots, [\lambda] + 1\}\}. \end{aligned}$$

As we have

$$\Delta_\mu(\Lambda^{-j} v(\Lambda^{-j} \cdot)) = \Lambda^{-j} (\Delta_{\mu \Lambda^j} v)(\Lambda^{-j} \cdot),$$

one gets, due to the scaling of the  $L^2$  norm, that

$$\mathcal{E}(v, \Lambda, \lambda) = \int_0^\infty \sum_{(j_1, j_2, j_3, j_4) \in \tilde{\mathcal{J}}} \|\Delta_{\mu \Lambda^{j_1}} v\|_{L^2} \|\Delta_{\mu \Lambda^{j_2}} v\|_{L^2} \|\Delta_{\mu \Lambda^{j_3}} v\|_{L^2} \|\Delta_{\mu \Lambda^{j_4}} v\|_{L^2} \frac{d\mu}{\mu}.$$

Without any loss of generality, we can assume that  $j_2$  is smaller than  $j_1$  in the above integral. Hölder inequality, for the measure  $\frac{d\mu}{\mu}$ , and the scaling invariance of this same measure, imply that

$$\begin{aligned} &\int_0^\infty \|\Delta_{\mu \Lambda^{j_1}} v\|_{L^2} \|\Delta_{\mu \Lambda^{j_2}} v\|_{L^2} \|\Delta_{\mu \Lambda^{j_3}} v\|_{L^2} \|\Delta_{\mu \Lambda^{j_4}} v\|_{L^2} \frac{d\mu}{\mu} \\ &\leq \|v\|_{\dot{B}_{2,4}^0(\mathbb{R}^2)}^2 \left( \int_0^\infty \|\Delta_\mu v\|_{L^2}^2 \|\Delta_{\mu \Lambda^{-(j_1-j_2)}} v\|_{L^2}^2 \frac{d\mu}{\mu} \right)^{\frac{1}{2}}. \end{aligned}$$

One deduces that:

$$\mathcal{E}(v, \Lambda, \lambda) \lesssim ([\lambda] + 1)^3 \|v\|_{\dot{B}_{2,4}^0(\mathbb{R}^2)}^2 \sum_{j=0}^{[\lambda]+1} \left( \int_0^\infty \|\Delta_\mu v\|_{L^2}^2 \|\Delta_{\mu \Lambda^{-j}} v\|_{L^2}^2 \frac{d\mu}{\mu} \right)^{\frac{1}{2}}.$$

The problem is reduced to the study of an integral of the form

$$\int_0^\infty f(\mu) g(\Lambda^{-1} \mu) \frac{d\mu}{\mu}.$$

with  $f$  and  $g$  in  $L^2(]0, \infty[; \mu^{-1}d\mu)$ . We define

$$\tilde{f}(\mu) \stackrel{\text{def}}{=} \mathbb{1}_{]0, \infty[}(\mu) \frac{f(\mu)}{\mu^{\frac{1}{2}}}.$$

Let us now write that

$$\int_0^\infty f(\mu)g(\Lambda^{-1}\mu)\frac{d\mu}{\mu} = \int_{\mathbb{R}} \tilde{f}(\mu)\Lambda^{-\frac{1}{2}}\tilde{g}(\Lambda^{-1}\mu) d\mu.$$

Let us apply Lemma 2.3 with  $d = 1$ ,  $p = 2$  and let us choose  $\varepsilon = C([\lambda] + 1)^{-4}$ ; this gives Inequality (6). In order to prove the entire proposition, let us concentrate on the continuity of  $\mathcal{F}$  which comes mostly from the following lemma.

**Lemma 3.3.** *The map*

$$\mathcal{D} \begin{cases} \dot{B}_{2,4}^0(\mathbb{R}^2) \times \mathbb{R}_+^* & \longrightarrow \dot{B}_{2,4}^0(\mathbb{R}^2) \\ (v, \delta) & \longmapsto \delta^{-1}v(\delta^{-1}\cdot) \end{cases}$$

is a continuous one.

*Proof.* It is worth noting that this map is an isometry in the sense that

$$\|\mathcal{D}(v, \delta)\|_{\dot{B}_{2,4}^0} = \|v\|_{\dot{B}_{2,4}^0}.$$

For any  $(v_0, \delta_0)$  and  $(v, \delta)$  belonging to  $\dot{B}_{2,4}^0(\mathbb{R}^2) \times \mathbb{R}_+^*$ ,

$$\begin{aligned} \|\mathcal{D}(v, \delta) - \mathcal{D}(v_0, \delta_0)\|_{\dot{B}_{2,4}^0(\mathbb{R}^2)} &\leq \|\delta^{-1}(v - v_0)(\delta^{-1}\cdot)\|_{\dot{B}_{2,4}^0(\mathbb{R}^2)} + \|\delta^{-1}v_0(\delta^{-1}\cdot) - \delta_0^{-1}v_0(\delta_0^{-1}\cdot)\|_{\dot{B}_{2,4}^0(\mathbb{R}^2)} \\ &\leq \|v - v_0\|_{\dot{B}_{2,4}^0(\mathbb{R}^2)} + \|v_0 - \delta\delta_0^{-1}v_0(\delta\delta_0^{-1}\cdot)\|_{\dot{B}_{2,4}^0(\mathbb{R}^2)}. \end{aligned}$$

Let us then consider a strictly positive real number  $\varepsilon$ . There exists a function  $v_{0,\varepsilon}$  of  $\mathcal{D}(\mathbb{R}^2)$  such that

$$\|v_0 - v_{0,\varepsilon}\|_{\dot{B}_{2,4}^0(\mathbb{R}^2)} \leq \frac{\varepsilon}{4}.$$

This yields then, for any strictly positive real number  $\delta$ ,

$$\begin{aligned} \|v_0 - \delta^{-1}v_0(\delta^{-1}\cdot)\|_{\dot{B}_{2,4}^0(\mathbb{R}^2)} &\leq \|v_0 - v_{0,\varepsilon}\|_{\dot{B}_{2,4}^0(\mathbb{R}^2)} + \|v_{0,\varepsilon} - \delta^{-1}v_{0,\varepsilon}(\delta^{-1}\cdot)\|_{\dot{B}_{2,4}^0(\mathbb{R}^2)} \\ &\quad + \|\delta^{-1}v_{0,\varepsilon}(\delta^{-1}\cdot) - \delta^{-1}v_0(\delta^{-1}\cdot)\|_{\dot{B}_{2,4}^0(\mathbb{R}^2)} \\ &\leq 2\|v_0 - v_{0,\varepsilon}\|_{\dot{B}_{2,4}^0(\mathbb{R}^2)} + \|v_{0,\varepsilon} - \delta^{-1}v_{0,\varepsilon}(\delta^{-1}\cdot)\|_{\dot{B}_{2,4}^0(\mathbb{R}^2)} \\ &\leq 2\frac{\varepsilon}{4} + C\|v_{0,\varepsilon} - \delta^{-1}v_{0,\varepsilon}(\delta^{-1}\cdot)\|_{L^2(\mathbb{R}^2)}. \end{aligned}$$

As  $\lim_{\delta \rightarrow 1} \|v_{0,\varepsilon} - \delta^{-1}v_{0,\varepsilon}(\delta^{-1}\cdot)\|_{L^2(\mathbb{R}^2)} = 0$ , one can infer that the mapping  $\mathcal{D}$  is a continuous one.  $\square$

*Conclusion of the proof of Proposition 2.1* The above lemma assures that the function  $\mathcal{F}$  is continuous on  $L^2(\mathbb{R}^2) \times \mathbb{R}_+^* \times (\mathbb{R}^+ \setminus \mathbb{N}^*)$ . In order to conclude the proof, let us observe that, for any positive integer  $k$ ,

$$\forall \lambda \in \left] -\frac{1}{4} + k, k + \frac{1}{4} \right[, \mathcal{F}(v, \Lambda, \lambda) = \sum_{j=0}^k \Lambda^{-j}v(\Lambda^{-j}\cdot) \quad (12)$$

which, thanks to Lemma 3.3, ensures the continuity of the function  $\mathcal{F}$  on  $\dot{B}_{2,4}^0(\mathbb{R}^2) \times \mathbb{R}_+^* \times \mathbb{R}^+$ . If  $\lambda$  belongs to the interval  $] -1/4 + k, k[$ , then  $[\lambda] = k - 1$ ,  $\chi(\lambda - [\lambda]) = 1$ , and Identity (12) is satisfied. If, now,  $\lambda$  belongs to the interval  $[k, k + 1/4[$ , then  $[\lambda] = k$ ,  $\chi(\lambda - [\lambda]) = 0$ , and the identity (12) is, again, satisfied. Proposition 2.1 is then proved.  $\square$



## 4 Almost-orthogonality and making of global solutions

We hereafter aim at proving Proposition 2.2. Classically, one searches the solution related to  $\mathcal{F}(v, \Lambda, \lambda)$  under the form  $u_{\text{app}} + R$ , where the approached solution  $u_{\text{app}}(t, x)$  is given by

$$\begin{aligned} u_{\text{app}}(t, x) &\stackrel{\text{def}}{=} \sum_{j=0}^{[\lambda]+1} V_{j,\Lambda}(t, x) \quad \text{with} \\ V_{j,\Lambda}(t, x) &\stackrel{\text{def}}{=} \Lambda^{-j} \mathcal{NLS}(v)(\Lambda^{-2j} t, \Lambda^{-j} x) \quad \text{for } j \text{ in } \{1, \dots, [\lambda]\} \quad \text{and} \\ V_{[\lambda]+1,\Lambda}(t, x) &\stackrel{\text{def}}{=} \Lambda^{-([\lambda]+1)} \mathcal{NLS}(\chi(\lambda - [\lambda])v)(\Lambda^{-2([\lambda]+1)} t, \Lambda^{-([\lambda]+1)} x). \end{aligned}$$

and where the term  $R$  has to be understood as an error one.

The idea is that if we choose the scaling parameter  $\Lambda$  large enough, then the error term  $R$  is small, due to the fact that solutions of very different scale almost do not interact. The point here is thus to prove that the choice of the size of the parameter  $\Lambda$  can be made continuously when the function  $v$  varies. In order to do so, let us write the equation satisfied by the error term  $R$ . One has

$$\begin{cases} i \partial_t R + \Delta R &= \tilde{E}_\Lambda + \Phi(R) \\ R|_{t=0} &= 0 \end{cases} \quad \text{with}$$

$$|\tilde{E}_\Lambda(t, x)| \lesssim E_\Lambda(t, x) \stackrel{\text{def}}{=} \sum_{\substack{0 \leq j, k, \ell \leq [\lambda]+1 \\ (j, k, \ell) \neq (j, j, j)}} |V_{j,\Lambda}(t, x)| |V_{k,\Lambda}(t, x)| |V_{\ell,\Lambda}(t, x)| \quad \text{and}$$

$$|\Phi(R)(t, x)| \lesssim |R(t, x)|^3 + |u_{\text{app}}(t, x)|^2 |R(t, x)|.$$

Here  $\Phi$  is a generic notation used to denote a function of  $R$  that will not be explicitly given, for the sake of simplicity.

Let us estimate  $\|E_\Lambda\|_{L^{\frac{4}{3}}(\mathbb{R}^{1+2})}$ . By definition of the functions  $V_{j,\Lambda}$ , we get using the Hölder inequality in the space variable

$$\begin{aligned} E_\Lambda^{j,k,\ell} &\stackrel{\text{def}}{=} \int_{\mathbb{R}^{1+2}} |V_{j,\Lambda}(t, x)|^{\frac{4}{3}} |V_{k,\Lambda}(t, x)|^{\frac{4}{3}} |V_{\ell,\Lambda}(t, x)|^{\frac{4}{3}} dx dt \\ &\leq \int_{\mathbb{R}} \Lambda^{-\frac{2(j+k+\ell)}{3}} \|\mathcal{NLS}(v)(\Lambda^{-2j} t, \cdot)\|_{L^4(\mathbb{R}^2)}^{\frac{4}{3}} \\ &\quad \times \|\mathcal{NLS}(v)(\Lambda^{-2k} t, \cdot)\|_{L^4(\mathbb{R}^2)}^{\frac{4}{3}} \|\mathcal{NLS}(v)(\Lambda^{-2\ell} t, \cdot)\|_{L^4(\mathbb{R}^2)}^{\frac{4}{3}} dt. \end{aligned}$$

As  $(j, k, \ell)$  is not of the type  $(j, j, j)$ , we can assume up to a perturbation of indices that  $k$  is smaller than  $\ell$ . Hölder inequality and then a change of variable imply that

$$\begin{aligned} E_\Lambda^{j,k,\ell} &\leq \left( \int_{\mathbb{R}} \|\mathcal{NLS}(v)(\Lambda^{-2j} t, \cdot)\|_{L^4(\mathbb{R}^2)}^4 \Lambda^{-2j} dt \right)^{\frac{1}{3}} \\ &\quad \times \left( \int_{\mathbb{R}} \|\mathcal{NLS}(v)(\Lambda^{-2k} t, \cdot)\|_{L^4(\mathbb{R}^2)}^2 \Lambda^{-(\ell-k)} \|\mathcal{NLS}(v)(\Lambda^{-2\ell} t, \cdot)\|_{L^4(\mathbb{R}^2)}^2 \Lambda^{-2k} dt \right)^{\frac{2}{3}} \\ &\leq \|\mathcal{NLS}(v)\|_{L^4(\mathbb{R}^{1+2})}^{\frac{4}{3}} \left( \int_{\mathbb{R}} \|\mathcal{NLS}(v)(t, \cdot)\|_{L^4(\mathbb{R}^2)}^2 \Lambda^{-(\ell-k)} \|\mathcal{NLS}(v)(\Lambda^{-2(\ell-k)} t, \cdot)\|_{L^4(\mathbb{R}^2)}^2 dt \right)^{\frac{2}{3}}. \end{aligned}$$

Applying Lemma 2.3, with  $d = 1$ ,  $p = 2$ ,  $f = g = \|\mathcal{NLS}(v)(t, \cdot)\|_{L^4(\mathbb{R}^2)}^2$ , one gets

$$\begin{aligned} \Lambda \geq \Lambda_F(\varepsilon, v) &\stackrel{\text{def}}{=} \Lambda_2(\varepsilon, \pi_{\mathbb{S}}^{(2)} \|\mathcal{NLS}(v)(t, \cdot)\|_{L^4(\mathbb{R}^2)}^2, \pi_{\mathbb{S}}^{(2)} \|\mathcal{NLS}(v)(t, \cdot)\|_{L^4(\mathbb{R}^2)}^2) \\ &\implies E_\Lambda^{j,k,\ell} \leq \varepsilon^{\frac{2}{3}} \|\mathcal{NLS}(v)\|_{L^4(\mathbb{R}^{1+2})}^4. \end{aligned}$$

As the function  $v$  is assumed to be in the ball  $B_\rho$ , we have  $\|\mathcal{NLS}(v)\|_{L^4(\mathbb{R}^{1+2})} \lesssim \|v\|_{\dot{B}_{2,4}^0(\mathbb{R}^2)}$ . Moreover the map

$$\begin{cases} B_\rho & \longrightarrow & L^2(\mathbb{R}) \\ v & \longmapsto & \|\mathcal{NLS}(v)(t, \cdot)\|_{L^4(\mathbb{R}^2)}^2 \end{cases}$$

is continuous. Thus the map  $(\varepsilon, v) \longmapsto \Lambda_F(\varepsilon, v)$  is continuous from  $]0, \infty[ \times B_\rho$  into  $]0, \infty[$ , and we have

$$\Lambda \geq \Lambda_F(\varepsilon, v) \implies E_\Lambda^{j,k,\ell} \lesssim \varepsilon^{\frac{2}{3}} \|v\|_{\dot{B}_{2,4}^0(\mathbb{R}^2)}^4. \quad (13)$$

By definition of  $E_\Lambda$  and  $E_\Lambda^{j,k,\ell}$ , we have

$$\|E_\Lambda\|_{L^{\frac{4}{3}}(\mathbb{R}^{1+2})}^{\frac{4}{3}} \lesssim \sum_{\substack{0 \leq j,k,\ell \leq [\lambda]+1 \\ (j,k,\ell) \neq (j,j,j)}} E_\Lambda^{j,k,\ell}.$$

With (13) this leads to

$$\Lambda \geq \Lambda_F(\varepsilon, v) \implies \|E_\Lambda\|_{L^{\frac{4}{3}}(\mathbb{R}^{1+2})} \leq C \varepsilon^{\frac{1}{2}} ([\lambda] + 1)^{\frac{9}{4}} \|v\|_{\dot{B}_{2,4}^0(\mathbb{R}^2)}^3. \quad (14)$$

The estimate of  $\|u_{\text{app}}\|_{L^4(\mathbb{R}^{1+2})}$  follows the same lines. Indeed, let us write that

$$\begin{aligned} & \left| \|u_{\text{app}}\|_{L^4(\mathbb{R}^{1+2})}^4 - \sum_{j=0}^{[\lambda]} \|V_{j,\Lambda}\|_{L^4(\mathbb{R}^{1+2})}^4 \right| \\ & \leq \sum_{\substack{0 \leq j_1, j_2, j_3, j_4 \leq [\lambda]+1 \\ j_1, j_2, j_3, j_4 \neq (j,j,j,j)}} \int_{\mathbb{R}^{1+2}} |V_{j_1,\Lambda}(t, x)| |V_{j_2,\Lambda}(t, x)| |V_{j_3,\Lambda}(t, x)| |V_{j_4,\Lambda}(t, x)| dx dt. \end{aligned}$$

Up to a permutation of indices, we can assume that  $j_1 \neq j_2$ . Hölder inequalities and scaling invariance lead to

$$\begin{aligned} & \left| \|u_{\text{app}}\|_{L^4(\mathbb{R}^{1+2})}^4 - \sum_{j=0}^{[\lambda]+1} \|V_{j,\Lambda}\|_{L^4(\mathbb{R}^{1+2})}^4 \right| \\ & \lesssim ([\lambda] + 1)^3 \|v\|_{\dot{B}_{2,4}^0(\mathbb{R}^2)}^2 \left( \sum_{j=1}^{[\lambda]+1} \int_{\mathbb{R}} \|\mathcal{NLS}(v)(t, \cdot)\|_{L^4(\mathbb{R}^2)}^2 \Lambda^{-j} \|\mathcal{NLS}(v)(\Lambda^{-2j}t, \cdot)\|_{L^4(\mathbb{R}^2)}^2 dt \right)^{\frac{1}{2}}. \end{aligned}$$

Applying Lemma 2.3, with  $d = 1$ ,  $p = 2$  and  $f = g = \|\mathcal{NLS}(v)(t, \cdot)\|_{L^4(\mathbb{R}^2)}^2$  gives

$$\Lambda \geq \Lambda_F(\varepsilon, v) \implies \left| \|u_{\text{app}}\|_{L^4(\mathbb{R}^{1+2})}^4 - ([\lambda] + 1) \|\mathcal{NLS}(v)\|_{L^4(\mathbb{R}^{1+2})}^4 \right| \leq C ([\lambda] + 1)^4 \varepsilon^{\frac{1}{2}} \|v\|_{\dot{B}_{2,4}^0(\mathbb{R}^2)}^4.$$

One deduces, then

$$\left( \varepsilon \leq \frac{1}{4([\lambda] + 1)^6} \quad \text{and} \quad \Lambda \geq \Lambda_F(\varepsilon, v) \right) \implies \|u_{\text{app}}\|_{L^4(\mathbb{R}^{1+2})}^4 \leq 2C ([\lambda] + 1) \|v\|_{\dot{B}_{2,4}^0(\mathbb{R}^2)}^4. \quad (15)$$

Let us admit for a while that, for any time  $T$  smaller than the maximal time of existence  $T^*$ ,

$$\|R\|_{L^4([0, T] \times \mathbb{R}^2)} \leq C (\|E_\Lambda\|_{L^{\frac{4}{3}}(\mathbb{R}^{1+2})} + \|R\|_{L^4([0, T] \times \mathbb{R}^2)}^2) \exp(\|u_{\text{app}}\|_{L^4(\mathbb{R}^{1+2})}^4). \quad (16)$$

A positive real number  $\eta$  being given, let us introduce the time  $T_\eta$  as

$$T_\eta \stackrel{\text{def}}{=} \sup \left\{ 0 \leq T < T^* / \int_0^T \|R(t, \cdot)\|_{L^4(\mathbb{R}^2)}^4 dt \leq \eta^4 \right\}.$$

If we prove that  $\bar{T} = T^*$  then the maximal time of existence  $T^*$  will be equal to  $+\infty$ . Using Inequality (16) we get that for any  $T$  smaller than  $T_\eta$ ,

$$\|R\|_{L^4([0,T]\times\mathbb{R}^2)} \leq C(\|E_\Lambda\|_{L^{\frac{4}{3}}(\mathbb{R}^{1+2})} + \eta^2) \exp(C\|u_{\text{app}}\|_{L^4(\mathbb{R}^{1+2})}^4).$$

Let us choose  $\eta$  and  $\varepsilon$  such that

$$\begin{aligned} \eta &= \frac{1}{4C} \exp(-C\|u_{\text{app}}\|_{L^4(\mathbb{R}^{1+2})}^4) \quad \text{and} \\ \varepsilon &= \frac{\eta^2}{4C^4([\lambda] + 1)^{\frac{9}{2}}\|v\|_{\dot{B}_{2,4}^0(\mathbb{R}^2)}^6} \exp(-4C([\lambda] + 1)\|v\|_{\dot{B}_{2,4}^0(\mathbb{R}^2)}^4). \end{aligned}$$

Assertions (14), (15) and (16) lead to  $\|R\|_{L^4([0,T]\times\mathbb{R}^2)} \leq 3/4$ . Thus,  $T_\eta = T^*$  which leads to  $T^* = +\infty$ . Our theorem is proved provided it is the case of inequality (16). Let us introduce the increasing sequence  $(T_m)_{0 \leq m \leq M+1}$  such that  $T_0 = 0$ ,  $T_{M+1} = +\infty$ , and

$$\forall m < M, \int_{T_m}^{T_{m+1}} \|u_{\text{app}}(t)\|_{L^4(\mathbb{R}^2)}^4 dt = c_0^4 \quad \text{and} \quad \int_{T_M}^{T_{M+1}} \|u_{\text{app}}(t)\|_{L^4(\mathbb{R}^2)}^4 dt \leq c_0^4 \quad (17)$$

for some given positive real number  $c_0$ , which will be chosen small enough later on. Obviously, we have

$$Mc_0^4 \leq \int_0^{T_M} \|u_{\text{app}}(t)\|_{L^4(\mathbb{R}^2)}^4 dt \leq \int_0^{+\infty} \|u_{\text{app}}(t)\|_{L^4(\mathbb{R}^2)}^4 dt \leq (M+1)c_0^4. \quad (18)$$

Thus the number  $M$  of  $T_m$  such that  $T_m$  is finite is less than  $c_0^{-4}\|u_{\text{app}}\|_{L^4(\mathbb{R}^{1+2})}^4$ . Let us set

$$\mathcal{R}(T) \stackrel{\text{def}}{=} \|R\|_{L^\infty([0,T];L^2(\mathbb{R}^2))} + \|R\|_{L^4([0,T]\times\mathbb{R}^2)}.$$

In order to prove, by induction, that for any  $T$  smaller than  $\min\{T_m, T^*\}$ , one has

$$\mathcal{R}(T) \leq C_0^{m+1} \|E_\Lambda\|_{L^{4/3}(\mathbb{R}^{1+2})}.$$

Let us now consider a time  $T$  smaller than  $\min\{T_{m+1}, T^*\}$ . Strichartz estimate to the interval  $[T_m, T]$  gives

$$\begin{aligned} \|R\|_{L^\infty([T_m,T];L^2(\mathbb{R}^2))} + \|R\|_{L^4([T_m,T]\times\mathbb{R}^2)} &\leq C(\|R(T_m)\|_{L^2(\mathbb{R}^2)} + \|E_\Lambda\|_{L^{\frac{4}{3}}(\mathbb{R}^{1+2})} \\ &\quad + \|R\|_{L^4([T_m,T]\times\mathbb{R}^2)}^3 + \|R\|_{L^4([T_m,T]\times\mathbb{R}^2)} \|u_{\text{app}}\|_{L^4([T_m,T_{m+1}]\times\mathbb{R}^2)}^2). \end{aligned}$$

By definition of  $c_0$ , we get

$$\begin{aligned} \|R\|_{L^\infty([T_m,T];L^2)} + \|R\|_{L^4([T_m,T]\times\mathbb{R}^2)} &\leq C(\|R(T_m)\|_{L^2(\mathbb{R}^2)} + \|E_\Lambda\|_{L^{\frac{4}{3}}(\mathbb{R}^{1+2})} \\ &\quad + (\eta^2 + c_0^2) \|R\|_{L^4([T_m,T]\times\mathbb{R}^2)}). \end{aligned}$$

If  $C(\eta^2 + c_0^2) \leq 1/2$ , one obtains

$$\|R\|_{L^\infty([T_m,T];L^2)} + \|R\|_{L^4([T_m,T]\times\mathbb{R}^2)} \leq 2C(\|R(T_m)\|_{L^2(\mathbb{R}^2)} + \|E_\Lambda\|_{L^{\frac{4}{3}}(\mathbb{R}^{1+2})}).$$

With the choice  $C_0 = 2C + 1$ , the induction hypothesis immediately yields

$$\begin{aligned} \mathcal{R}(T) &\leq (2C + 1)C_0^{m+1} \|E_\Lambda\|_{L^{\frac{4}{3}}(\mathbb{R}^{1+2})} \\ &\leq C_0^{m+2} \|E_\Lambda\|_{L^{\frac{4}{3}}(\mathbb{R}^{1+2})}. \end{aligned}$$

By induction, we deduce that, for any  $T$  less than  $T^*$ , we have  $\mathcal{R}(T) \leq C_0^{M+2} \|E_\Lambda\|_{L^{\frac{4}{3}}(\mathbb{R}^{1+2})}$ . Using Inequality (18), it turns out that

$$\mathcal{R}(T) \leq C \|E_\Lambda\|_{L^{\frac{4}{3}}(\mathbb{R}^{1+2})} \exp\left(C \int_0^\infty \|u_{\text{app}}(t)\|_{L^4(\mathbb{R}^2)}^4 dt\right)$$

which imply Inequality (16). Proposition 2.2 is proved.  $\square$

## 5 Some non-increasing reordering properties

We hereafter aim at proving proposition 3.1. More information can be found in [9], [16]. We nonetheless recall the main results useful to our study. Various proofs of those results can be found in the above references. In the following, we concentrate on positive-valued functions. We will frequently refer to the fact that, for any function  $h$ ,

$$h(x) = \int_0^\infty \mathbb{1}_{(h>t)}(x) dt. \quad (19)$$

To prove (7), let us note that  $(f(s \cdot) > t) = s^{-1}(f > t)$ . Let  $R_t$  denote the radius such that the volume of  $(f > t)$  is the same as the one of the centered ball of radius  $R_t$ . Thus we have

$$\mathbb{1}_{(f(s \cdot) > t)^*}(x) = \mathbb{1}_{B(0, s^{-1}R_t)}(x) = \mathbb{1}_{(f > t)^*}(sx).$$

One obtains inequality (7) by integration in  $t$ .

The proof of Inequality (8) relies on the following observation. Let  $A$  and  $B$  denote two sets of  $\mathbb{R}^d$  of finite measure. If  $A^*$  and  $B^*$  are two centered balls, one has

$$\text{meas}(A^* \cap B^*) = \min\{\text{meas } A^*, \text{meas } B^*\} = \min\{\text{meas } A, \text{meas } B\} \geq \text{meas}(A \cap B). \quad (20)$$

This can be illustrated by the following drawing.

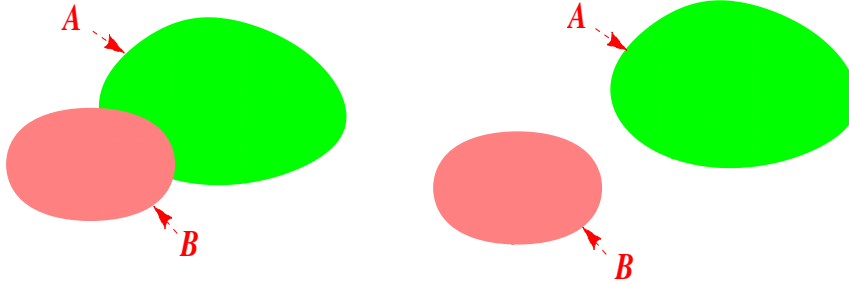


Figure 3: The sets  $A$  et  $B$ , with or without an intersection.

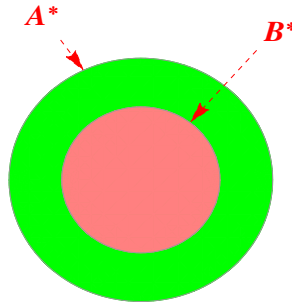


Figure 4: The centered balls  $A^*$  and  $B^*$ .

Thanks to Fubini's theorem, and to inequality (20), one has, from the definition of  $f^*$  and  $g^*$ ,

$$\begin{aligned} \int_{\mathbb{R}^d} f^*(x)g^*(x)dx &= \int_{\mathbb{R}^d} \left( \int_0^\infty \mathbb{1}_{(f>t)^*}(x) dt \right) \left( \int_0^\infty \mathbb{1}_{(g>s)^*}(x) ds \right) dx \\ &= \int_0^\infty \int_0^\infty \text{meas}((f > t)^* \cap (g > s)^*) dt ds. \end{aligned}$$

By applying (20) in the above relation, one gets:

$$\int_{\mathbb{R}^d} f^*(x)g^*(x)dx \geq \int_0^\infty \int_0^\infty \text{meas}((f > t) \cap (g > s)) dt ds.$$

If we apply then (19), we obtain:

$$\begin{aligned} \int_{\mathbb{R}^d} f^*(x)g^*(x)dx &\geq \int_0^\infty \int_0^\infty \int_{\mathbb{R}^d} \mathbb{1}_{(f>t)}(x)\mathbb{1}_{(g>s)}(x) dt ds dx \\ &\geq \int_{\mathbb{R}^d} \left( \int_0^\infty \mathbb{1}_{(f>t)}(x) dt \right) \left( \int_0^\infty \mathbb{1}_{(g>s)}(x) ds \right) dx \\ &\geq \int_{\mathbb{R}^d} f(x)g(x) dx \end{aligned}$$

which ensures the required inequality.

The proof of Inequality (9) is a very close one. According to (19) and Fubini's theorem, we get

$$\begin{aligned} \int_{\mathbb{R}^d} f(x)\mathbb{1}_{(g\leq s)}(x)dx &= \int_{\mathbb{R}^d} \int_0^\infty \mathbb{1}_{(f>t)}(x)\mathbb{1}_{(g\leq s)}(x) dx dt \\ &= \int_0^\infty \text{meas}((f > t) \cap (g \leq s)) dt. \end{aligned} \quad (21)$$

As  $\mathbb{1}_{(f>t)}\mathbb{1}_{(g\leq s)} = \mathbb{1}_{(f>t)} - \mathbb{1}_{(f>t)}\mathbb{1}_{(g>s)}$ , an integration leads to:

$$\text{meas}((f > t) \cap (g \leq s)) = \text{meas}(f > t) - \text{meas}((f > t) \cap (g > s)).$$

By applying (20) to the sets  $(f > t)$  and  $(g > s)$ , one deduces

$$\begin{aligned} \text{meas}((f > t) \cap (g \leq s)) &\geq \text{meas}((f > t)^*) - \text{meas}((f > t)^* \cap (g > s)^*) \\ &\geq \text{meas}(f > t)^* - \text{meas}((f^* > t) \cap (g^* > s)) \\ &\geq \text{meas}((f > t)^* \cap (g^* \leq s)). \end{aligned}$$

Using it in inequality (21) enables one to write, thanks to Fubini's theorem, that

$$\begin{aligned} \int_{\mathbb{R}^d} f(x)\mathbb{1}_{(g\leq s)}(x)dx &\geq \int_0^\infty \int_{\mathbb{R}^d} \mathbb{1}_{(f>t)^*}(x)\mathbb{1}_{(g^*\leq s)}(x) dx dt \\ &\geq \int_{\mathbb{R}^d} f^*(x)\mathbb{1}_{(g^*\leq s)}(x) dx \end{aligned}$$

which is Inequality (9).

To prove the equality of the  $L^p$  norms, let us note that, according to Fubini's theorem, for any positive valued function  $h$ , and any  $p$  in  $[1, \infty[$ , one has

$$\begin{aligned} \|h\|_{L^p}^p &= p \int_{\mathbb{R}^d} \int_0^\infty \mathbb{1}_{(h>t)}(x)t^{p-1} dt dx \\ &= p \int_0^\infty \text{meas}(h > t)t^{p-1} dt. \end{aligned}$$

By definition of  $\star$  operation on sets,  $(f > t)$  and  $(f^* > t)$  are of same measure. The above formula ensures immediately the equality of the  $L^p$  norms.

The fact that the  $\star$  process is 1-lipschitzian on  $L^p(\mathbb{R}^d)$  comes from the more general following result.

**Lemma 5.1.** *Let  $J$  denote a convex function from  $\mathbb{R}$  to  $\mathbb{R}^+$  such that  $J(0) = 0$ . If  $f$  and  $g$  are two positive-valued functions, then*

$$\int_{\mathbb{R}^d} J(f^*(x) - g^*(x)) dx \leq \int_{\mathbb{R}^d} J(f(x) - g(x)) dx.$$

*Proof.* Let us write

$$J = J_+ + J_- \quad \text{with} \quad J_+ \stackrel{\text{def}}{=} \mathbb{1}_{\mathbb{R}^+} J \quad \text{and} \quad J_- \stackrel{\text{def}}{=} \mathbb{1}_{\mathbb{R}^- \setminus \{0\}} J.$$

The two functions  $J_+$  and  $J_-$  are convex, and positive valued. One proves the inequality for  $J_+$ , the proof for  $J_-$  being strictly analogous. Let  $h_1$  and  $h_2$  denote two non negative functions. One has

$$\begin{aligned} J_+(h_1(x) - h_2(x)) &= \int_{h_2(x)}^{h_1(x)} J'_+(h_1(x) - s) ds \\ &= \int_{h_2(x)}^{+\infty} J'_+(h_1(x) - s) \mathbb{1}_{(h_2 \leq s)}(x) ds. \end{aligned} \quad (22)$$

If  $H$  is an increasing function on  $\mathbb{R}^+$ , its derivative in the sense of distributions is a non negative measure (which we denote by  $dH$ ) and which satisfies

$$H(y - s) = \int_0^\infty \mathbb{1}_{]s+t, \infty[}(y) dH(t).$$

Thanks to Fubini's theorem, we deduce that, for any set of functions  $(h_1, h_2)$ ,

$$\int_{\mathbb{R}^d} H(h_1(x) - s) \mathbb{1}_{(h_2 \leq s)}(x) dx = \int_0^\infty \left( \int_{\mathbb{R}^d} \mathbb{1}_{(h_1 > s+t)}(x) \mathbb{1}_{(h_2 \leq s)}(x) dx \right) dH(t). \quad (23)$$

By applying this formula with  $(h_1, h_2) = (f^*, g^*)$ , in conjunction with Inequality (9), one obtains, thanks to the positivity of the measure  $dH(t)$ ,

$$\begin{aligned} \int_{\mathbb{R}^d} H(f^*(x) - s) \mathbb{1}_{(g^* \leq s)}(x) dx &= \int_{\mathbb{R}^d \times [0, \infty[} \mathbb{1}_{(f^* > s+t)}(x) \mathbb{1}_{(g^* \leq s)}(x) dx dH(t) \\ &\leq \int_0^\infty \left( \int_{\mathbb{R}^d} \mathbb{1}_{(f > s+t)}(x) \mathbb{1}_{(g \leq s)}(x) dx \right) dH(t). \end{aligned}$$

If, now, we apply Formula (23) with the set of functions  $(f, g)$ , we get

$$\int_{\mathbb{R}^d} H(f^*(x) - s) \mathbb{1}_{(g^* \leq s)}(x) dx \leq \int_{\mathbb{R}^d} H(f(x) - s) \mathbb{1}_{(g \leq s)}(x) dx.$$

By applying this inequality with  $H = J'_+$ , it ensures the required result, thanks to Formula (22).  $\square$

## References

- [1] H. Bahouri and P. Gérard, High frequency approximation of solutions to critical nonlinear wave equations, *American Journal of Math*, **121**, 1999, pages 131-175.
- [2] J. Bourgain, Refinements of Strichartz' inequality and applications to 2D-NLS with critical nonlinearity, *International Mathematical Research Notices*, **5**, 1998, pages 253-283.
- [3] N. Burq, Explosion pour l'équation de Schrödinger au régime du "log-log", (d'après Merle-Raphael), *Séminaire Bourbaki* 2005/06), *Astérisque* **311**, 2007, pages 33-53.
- [4] T. Cazenave and F. Weissler, Some remarks on the nonlinear Schrödinger equation in the critical case, *Nonlinear Semigroups, Partial Differential Equations and Attractors* Lecture Notes in Mathematics, **1394**, Berlin, 1989, pages 18-29.
- [5] (avec I. Gallagher) Wellposedness and stability results for the Navier-Stokes equations in  $\mathbf{R}^3$ , *Annales de l'Institut Henri Poincaré, Analyse Non Linéaire*, **26**, 2009, pages 599-624.

- [6] J.-Y. Chemin, I. Gallagher and P. Zhang, Sums of large global solutions to the incompressible Navier-Stokes equations, *Journal für reine und angewandte Mathematik*, **681**, 2013, pages 65-82.
- [7] B. Dodson, Global well-posedness and scattering for the defocusing,  $L^2$ -critical, nonlinear Schrödinger equation when  $d = 2$ , 2010, arXiv :1006.1375.
- [8] B. Dodson, Global well-posedness and scattering for the mass critical nonlinear Schrödinger equation with mass below the mass of the ground state, arXiv :1104.1114, 2011.
- [9] E. Lieb and M. Loss, Analysis, Graduate Studies in Mathematics, **14** (American Mathematical Society, Providence, RI, 1997).
- [10] F. Merle and P. Raphaël, On universality of blow up profile for  $L^2$  critical nonlinear Schrödinger equation, *Inventiones Mathematicae*, **156**, 2004, pages 565-672.
- [11] F. Merle, P. Raphaël, J. Szeftel, The instability of Bourgain-Wang solutions for  $L^2$  critical NSL, *American Journal of Mathematics*, **135**, 2013 pages 967-1017.
- [12] F. Merle and L. Vega, Compactness at blow-up time for  $L^2$  solutions of the critical nonlinear Schrödinger equation in 2D, *International Mathematical Research Notices*, **1998**, pages 399-425.
- [13] F. Planchon, Dispersive estimates and the 2D Cubic NLS equation, *Journal d'Analyse Mathématique*, **86**, 2002, pages 319-334.
- [14] F. Planchon, On the Cauchy problem in Besov spaces for a non-linear Schrödinger equation, *Communications in Contemporary Mathematics*, **2**, 2000, pages 243-254.
- [15] F. Planchon, Existence globale et scattering pour les solutions de masse finie de l'équation de Schrödinger cubique en dimension 2, Séminaire Bourbaki, 63<sup>ème</sup> année, 2010-2011, n° 1042.
- [16] G. Pólya and G. Szegő, *Isoperimetric inequalities in Mathematical Physics*, Annals of Mathematics Studies, Vol. 27, Princeton University Press, 1951.