

*On the Control of some Galpern-Sobolev
Equations*

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Outline

- Introduction
 - Barenblatt-Zhel'tov-Kochina equation
 - Benjamin-Bona-Mahony equation
- Controllability
 - Fixed control domain: lack of null controllability
 - A remedy
- Final Comments

Sobolev-Galpern equations

A partial differential equation is of Sobolev-Galpern type if the highest order terms contain derivatives in both space and time coordinates.

A typical example is

$$\mathcal{M}\partial_t u + \mathcal{L}u = f$$

with \mathcal{L} and \mathcal{M} are linear partial differential operators in the spatial variable of order $2m$ and $l \leq 2m$, respectively (independent of t).

- seepage of fluids through fissured rocks;
- unsteady flows of second-order fluids;
- consolidation;
- surface waves of long wavelength in liquids
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We will focus on the following equations:

- *Barenblatt-Zhel'tov-Kochina equation*

$$u_t - \Delta u_t - \Delta u = f$$

- *Linearized Benjamin-Bona-Mahony equation*

$$u_t - u_{xxt} + u_x = f$$

Barenblatt-Zhel'tov-Kochina equation

We are interested in the null controllability of the BZK equation:

$$\left\{ \begin{array}{ll} u_t - \Delta u_t - \Delta u = f1_\omega & \text{in } Q := (0, T) \times \Omega, \\ u = 0 & \text{on } \Sigma := (0, T) \times \partial\Omega, \\ u(0) = u_0 & \text{in } \Omega, \end{array} \right.$$

i.e, given u_0 , find f such that the solution of BZK satisfies:

$$u(T) = 0.$$

Here, $\Omega \subset \mathbb{R}^N$ and $\omega \subset \Omega$ is a nonempty subset.

Some properties of BZK

$$\left| \begin{array}{ll} u_t - \Delta u_t - \Delta u = 0 & \text{in } Q := (0, T) \times \Omega, \\ u = 0 & \text{on } \Sigma := (0, T) \times \partial\Omega, \\ u(0) = u_0 & \text{in } \Omega, \end{array} \right.$$

- In space, the solution is just as smooth as the initial data allow it to be;
- The solution is very regular in time;
- Null controllability implies exact controllability.

A nice decomposition

The BZK equation can be decomposed as

$$\left\{ \begin{array}{ll} u - \Delta u = v & \text{in } Q, \\ v_t + v - u = f1_\omega & \text{in } Q, \\ u = 0 & \text{on } \Sigma, \\ v(0) = u_0 - \Delta u_0 & \text{in } \Omega. \end{array} \right.$$

We have:

$$v(T) = 0 \text{ iff } u(T) = 0.$$

Observability inequality

The null controllability for BZK system is equivalent to the observability inequality for the adjoint solutions, i.e

$$\|\psi(\cdot, 0)\|_{L^2(\Omega)}^2 \leq C \iint_{\omega \times (0, T)} |\psi|^2,$$

where

$$\left\{ \begin{array}{ll} \varphi - \Delta \varphi = \psi & \text{in } Q, \\ -\psi_t + \psi = \varphi & \text{in } Q, \\ \varphi = 0 & \text{on } \Sigma, \\ \psi(T) = \psi_T & \text{in } \Omega. \end{array} \right.$$

Lack of controllability with fixed control

Remark

For a given (fixed) $\omega \subset \Omega$, null controllability fails for system BZK.

This is due to:

- the spectrum $(I - \Delta)^{-1} \Delta$ is given by $\{\frac{\mu_j}{1+\mu_j} : \mu_j \in \sigma(-\Delta)\}$;
- highly localized solutions (Gaussian beam);

What can we do?

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What can we do?

Remedy: Moving Control

Make the control in the second equation of BZK move on time, i.e. $\omega = \omega(t)$.

The set $\omega(t)$ covers the whole domain Ω in its motion.

This strategy has been previously used for other models:

L. Rosier, B.-Y. Zhang, Unique continuation property and control for the Benjamin-Bona-Mahony equation on a periodic domain, *J. Differential Equations* 254 (2013), 141–178.

P. Martin, L. Rosier, P. Rouchon, Null Controllability of the Structurally Damped Wave Equation with Moving Control, *SIAM J. Control Optim.*, 51 (1) (2013), 660–684.

F. W. Chaves-Silva, L. Rosier, and E. Zuazua, Null Controllability of a System of Viscoelasticity with Moving Control, *J. Math. Pures Appl.*, 101 (9) (2014), 198–222.

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BZK with Moving Control

With moving controls, BZK equation and BZK system read

$$\left\{ \begin{array}{ll} u_t - \Delta u - \Delta u_t = f1_{\omega(t)} & \text{in } Q, \\ u = 0 & \text{on } \Sigma, \\ u(\cdot, 0) = u_0 & \text{in } \Omega \end{array} \right.$$

and

$$\left\{ \begin{array}{ll} u - \Delta u = v & \text{in } Q, \\ v_t + v - u = f1_{\omega(t)} & \text{in } Q, \\ u = 0 & \text{on } \Sigma, \\ v(\cdot, 0) = v_0 & \text{in } \Omega, \end{array} \right.$$

respectively.

And, still, we want to find a control f such that $u(T) = v(T) = 0$.

1D case

For the 1D case, using moment method, Q. Tao et al.¹, showed the null controllability of BZK with periodic boundary condition.

¹Q. Tao, H. Gao, Z. Yao, Null controllability of a pseudo-parabolic equation with moving control, J. Math. Analysis and Appl., 418 (2)(2014)

N-dimensional case

The adjoint system of BZK reads

$$\left\{ \begin{array}{ll} \varphi - \Delta\varphi = \psi & \text{in } Q, \\ -\psi_t + \psi = \varphi & \text{in } Q, \\ \varphi = 0 & \text{on } \Sigma, \\ \psi(T) = \psi_T & \text{in } \Omega. \end{array} \right.$$

Null controllability of BZK system is equivalent to

$$\|\psi(\cdot, 0)\|_{L^2(\Omega)}^2 \leq C \int_0^T \int_{\omega(t)} |\psi|^2 dxdt.$$

Carleman inequality

Theorem (Chaves-Silva & S.)

Given $\psi_T \in L^2(\Omega)$, the solution (φ, ψ) of the adj. system of BZK satisfies:

$$\begin{aligned} & \iint_Q \rho_1(x, t)(|\nabla \varphi|^2 + |\varphi|^2) dx dt + \iint_Q \rho_2(x, t)|\psi|^2 dx dt \\ & + \iint_Q \rho_3(t)(|\nabla \varphi_t|^2 + |\varphi_t|^2) dx dt \\ & \leq C \int_0^T \int_{\omega(t)} \rho_4(x, t)|\psi|^2 dx dt, \end{aligned}$$

where ρ_i , $i = 1, \dots, 4$ are appropriate weights.

Idea of the proof

Three main difficulties appear:

- 1 Carleman inequalities for the Laplace operator and ODE equations with a moving control region²;
- 2 We must have **the same** weight functions in the Carleman for both equations.
- 3 Eliminate a local integral of φ .

Fortunately, we can handle all these difficulties.

²F. W. Chaves-Silva, L. Rosier, and E. Zuazua, Null controllability of a system of viscoelasticity with a moving control, J. Math. Pures Appl., 101 (9) (2014), 198–222.

First controllability result

Theorem (Chaves-Silva & S.)

Given $v_0 \in L^2(\Omega)$, there exists $f \in L^2(Q)$ such that the associated solution (u, v) of BZK system satisfies:

$$v(T) = u(T) = 0.$$

Null controllability for BZK

Corollary

Given $u_0 \in H_0^1(\Omega) \cap H^2(\Omega)$, there exists $f \in L^2(Q)$ such that the associated solution u of BZK equation satisfies

$$u(T) = 0.$$

Linearized Benjamin-Bona-Mahony equation

Similar ideas (but not the same!) can be used to study the controllability of

$$\left\{ \begin{array}{ll} u_t - \Delta u_t - \operatorname{div}(A(x, t)u) = f1_\omega & \text{in } Q, \\ u = 0 & \text{on } \Sigma, \\ u(0) = u_0 & \text{in } \Omega, \end{array} \right.$$

where A is a regular enough vector function.

A nice decomposition

The BBM equation can be decomposed as

$$\left\{ \begin{array}{ll} u - \Delta u = v & \text{in } Q, \\ v_t + \nabla \cdot (A(x, t)u) = f1_\omega & \text{in } Q, \\ u = 0 & \text{on } \Sigma, \\ v(\cdot, 0) = v_0 & \text{in } \Omega. \end{array} \right.$$

We have:

$$v(T) = 0 \text{ iff } u(T) = 0.$$

Observability inequality

The null controllability for BBM system is equivalent to the observability inequality for the adjoint solutions, i.e

$$\|\psi(\cdot, 0)\|_{L^2(\Omega)}^2 \leq C \iint_{\omega \times (0, T)} |\psi|^2,$$

where

$$\left\{ \begin{array}{ll} \varphi - \Delta \varphi = A \cdot \nabla \psi & \text{in } Q, \\ -\psi_t = \varphi & \text{in } Q, \\ \varphi = 0 & \text{on } \Sigma, \\ \psi(T) = \psi_T & \text{in } \Omega. \end{array} \right.$$

- Lack of controllability^{3 4}: easy to see when $\text{supp}(A) \cap \omega = \emptyset$;

³S. Micu, On the controllability of the linearized Benjamin-Bona-Mahony equation, *SIAM J. Control Optim.*, **39** (2001), 1677–1696.

⁴X. Zhang, E. Zuazua, Unique continuation for the linearized Benjamin-Bona-Mahony equation with space-dependent potential, *Mathematische Annalen* 325 (2003), 543–582.

Old same remedy: moving controls

With moving controls, BBM system reads

$$\left\{ \begin{array}{ll} u - \Delta u = v & \text{in } Q, \\ v_t + \nabla \cdot (A(x, t)u) = f1_{\omega(t)} & \text{in } Q, \\ u = 0 & \text{on } \Sigma, \\ v(\cdot, 0) = v_0 & \text{in } \Omega. \end{array} \right.$$

And, still, we want to find a control f such that $u(T) = v(T) = 0$.

The adjoint system of BBM reads

$$\left\{ \begin{array}{ll} \varphi - \Delta\varphi = A \cdot \nabla\psi & \text{in } Q, \\ -\psi_t = \varphi & \text{in } Q, \\ \varphi = 0 & \text{on } \Sigma, \\ \psi(T) = \psi_T & \text{in } \Omega. \end{array} \right.$$

Null controllability of BZK system is equivalent to

$$\|\psi(\cdot, 0)\|_{L^2(\Omega)}^2 \leq C \int_0^T \int_{\omega(t)} |\psi|^2 dxdt.$$

Theorem (Chaves-Silva & S.)

Given $\psi_T \in L^2(\Omega)$, the solution (φ, ψ) of the adj. system of BZK satisfies:

$$\begin{aligned} & \iint_Q \rho_1^*(x, t)(|\nabla \varphi|^2 + |\varphi|^2) dxdt + \iint_Q \rho_2^*(x, t)|\psi|^2 dxdt \\ & + \iint_Q \rho_3^*(t)(|\nabla \varphi_t|^2 + |\varphi_t|^2) dxdt \\ & \leq C \int_0^T \int_{\omega(t)} \rho_4^*(x, t)|\psi|^2 dxdt, \end{aligned}$$

where ρ_i^* , $i = 1, \dots, 4$ are appropriate weights.

Key ingredients:

- $A \cdot \nabla \psi = \nabla \cdot (A\psi) - \psi \nabla \cdot A$;
- H^{-1} elliptic Carleman inequality;
- energy estimates;

Final comments

- Moving control is a good alternative;
- Neumann boundary conditions:
 - 1D is OK!
 - N-dimensional?
- Controllability results for general Sobolev-Galpern type equations

$$\mathcal{M}\partial_t u + \mathcal{L}u = f1_\omega?$$

Happy birthday Jean-Michel!!!